

ECE 546

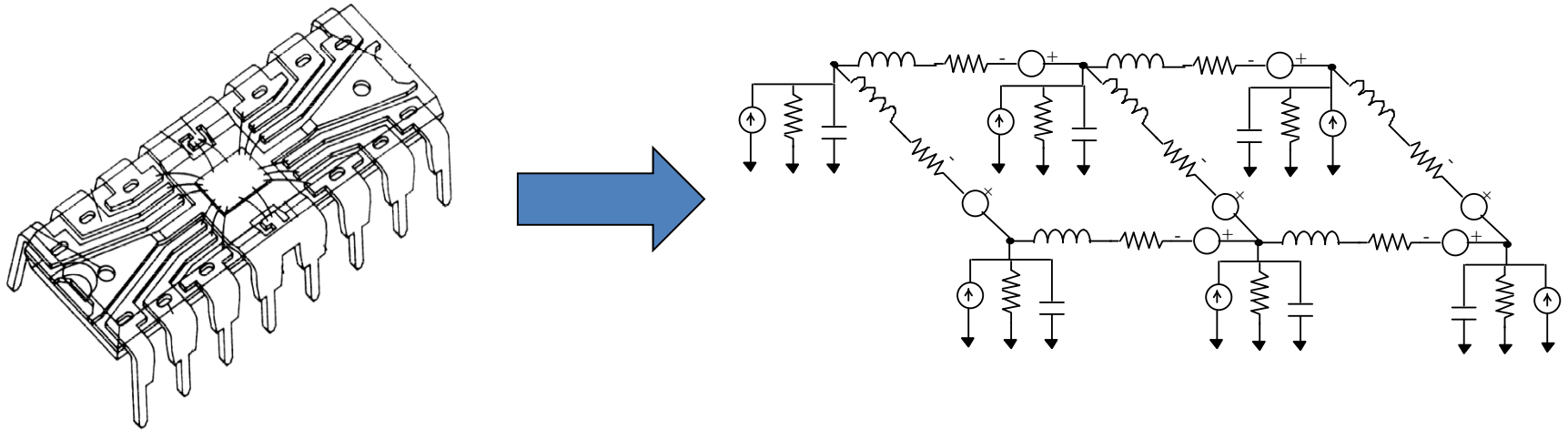
Lecture 04

Resistance, Capacitance, Inductance

Spring 2016

Jose E. Schutt-Aine
Electrical & Computer Engineering
University of Illinois
jschutt@emlab.uiuc.edu

What is Extraction?



Process in which a complex arrangement of conductors and dielectric is converted into a netlist of elements in a form that is amenable to circuit simulation.

Need Field Solvers

Electromagnetic Modeling Tools

We need electromagnetic modeling tools to analyze:

Transmission line propagation
Reflections from discontinuities
Crosstalk between interconnects
Simultaneous switching noise

So we can provide:

Improved design of interconnects
Robust design guidelines
Faster, more cost effective design cycles

Field Solvers – History

◆1960s

Conformal mapping techniques
Finite difference methods (2-D Laplace eq.)
Variational methods

◆1970s

Boundary element method
Finite element method (2-D)
Partial element equivalent circuit (3-D)

◆1980s

Time domain methods (3-D)
Finite element method (3-D)
Moment method (3-D)
rPEEC method (3-D)

◆1990s

Adapting methods to parallel computers
Including methods in CAD tools

◆2000s

Incorporation of Passivity
Incorporation of Causality

◆2010s

Stochastic Techniques
Multiphysics Tools

Categories of Field Solvers

- Method of Moments (MOM)
- Application to 2-D Interconnects
- Closed-Form Green's Function
- Full-Wave and FDTD
- Parallel FDTD
- Applications

Capacitance

Relation: $Q = Cv$

Q : charge stored by capacitor

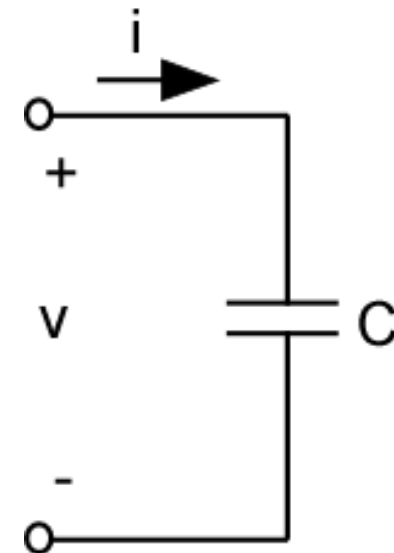
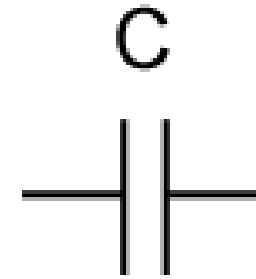
C : capacitance

v : voltage across capacitor

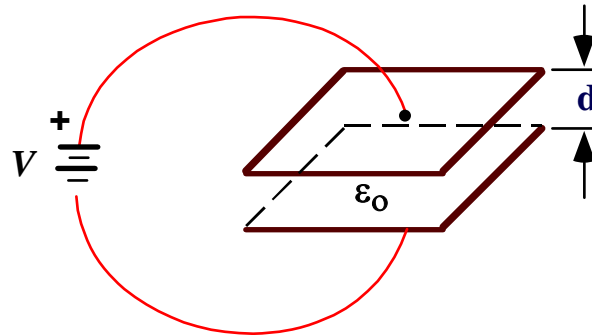
i : current into capacitor

$$i(t) = C \frac{dv}{dt} = \frac{dQ}{dt}$$

$$v(t) = \frac{1}{C} \int_0^t i(\tau) d\tau$$



Capacitance



$$C = \frac{\epsilon_0 A}{d}$$

A : area

ε₀ : permittivity

For more complex capacitance geometries, need to use numerical methods

Capacitance Calculation

$$\phi(r) = \int g(r, r') \sigma(r') dr'$$

$\phi(r)$ = potential (*known*)

$g(r, r')$ = Green's function (*known*)

$\sigma(r')$ = charge distribution (*unknown*)

Once the charge distribution is known, the total charge Q can be determined. If the potential $\phi=V$, we have

$$Q=CV$$

To determine the charge distribution, use the moment method

METHOD OF MOMENTS

Operator equation

$$L(\mathbf{f}) = \mathbf{g}$$

L = integral or differential operator

\mathbf{f} = unknown function

\mathbf{g} = known function

Expand unknown function \mathbf{f}

$$\mathbf{f} = \sum_n \alpha_n \mathbf{f}_n$$

METHOD OF MOMENTS

in terms of basis functions f_n , with unknown coefficients α_n to get

$$\sum_n \alpha_n L(f_n) = g$$

Finally, take the scalar or *inner product* with testing or weighting functions w_m :

$$\sum_n \alpha_n \langle w_m, Lf_n \rangle = \langle w_m, g \rangle$$

Matrix equation

$$[l_{mn}][\alpha_n] = [g_m]$$

METHOD OF MOMENTS

$$[l_{mn}] = \begin{bmatrix} \langle w_1, Lf_1 \rangle & \langle w_1, Lf_2 \rangle & \dots \\ \langle w_2, Lf_1 \rangle & \langle w_2, Lf_2 \rangle & \dots \\ \dots & \dots & \dots \end{bmatrix}$$

$$[\alpha_n] = \begin{bmatrix} \alpha_1 \\ \alpha_2 \\ \cdot \end{bmatrix}$$

$$[g_m] = \begin{bmatrix} \langle w_1, g \rangle \\ \langle w_2, g \rangle \\ \cdot \end{bmatrix}$$

Solution for weight coefficients

$$[\alpha_n] = [l_{nm}^{-1}] [g_m]$$

Moment Method Solution

$$\nabla \cdot D = \rho$$

$$E = -\nabla \phi$$

$$\nabla^2 \phi = -\frac{\rho}{\epsilon}$$

$$L\phi = -\frac{\rho}{\epsilon}$$

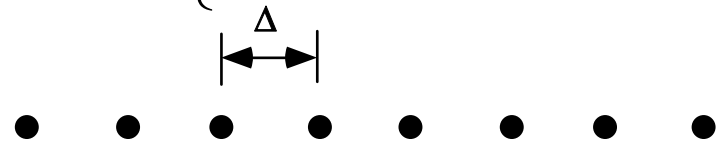
$$\phi(x, y, z) = \iiint \frac{\rho(x', y', z')}{4\pi\epsilon R} dx' dy' dz'$$

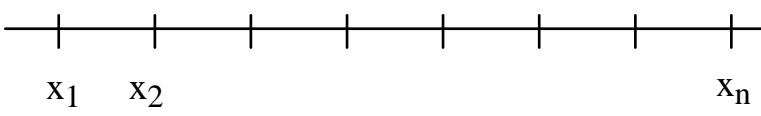
$$R = \sqrt{(x-x')^2 + (y-y')^2 + (z-z')^2}$$

Green's function G: $LG = \delta$

Basis Functions

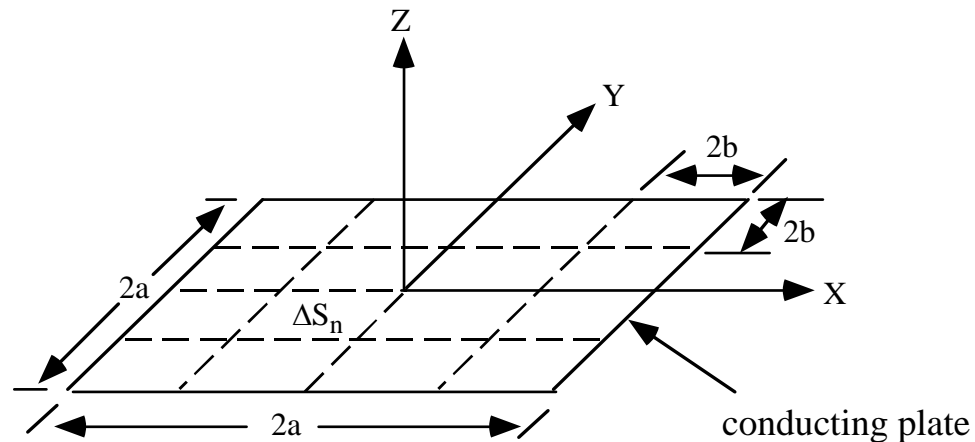
Subdomain bases

$$P(x_n) = \begin{cases} 1 & x_n - \frac{\Delta}{2} < x < x_n + \frac{\Delta}{2} \\ 0 & \textit{otherwise} \end{cases}$$


$$T(x_n) = \begin{cases} 1 - |x| & x_n - \frac{\Delta}{2} < x < x_n + \frac{\Delta}{2} \\ 0 & \textit{otherwise} \end{cases}$$


Testing functions often (not always) chosen same as basis function.

Conducting Plate



$$\phi(x, y, z) = \int_{-a}^a dx' \int_{-a}^a dy' \frac{\sigma(x', y', z')}{4\pi\epsilon R}$$

σ = charge density on plate

Conducting Plate

Setting $\phi = V$ on plate

$$R = \sqrt{(x - x')^2 + (y - y')^2}$$

$$V = \int_{-a}^a dx' \int_{-a}^a dy' \frac{\sigma(x', y', z')}{4\pi\epsilon \sqrt{(x - x')^2 + (y - y')^2}}$$

for $|x| < a$; $|y| < a$

Capacitance of plate: $C = \frac{q}{V} = \frac{1}{V} \int_{-a}^a dx' \int_{-a}^a dy' \sigma(x', y', z')$

Conducting Plate

Basis function P_n

$$P_n(x_m, y_n) = \begin{cases} 1 & x_m - \frac{\Delta s}{2} < x < x_m + \frac{\Delta s}{2} \\ 1 & y_n - \frac{\Delta s}{2} < y < y_n + \frac{\Delta s}{2} \\ 0 & \textit{otherwise} \end{cases}$$

Representation of unknown charge

$$\sigma(x, y) = \sum_{n=1}^N \alpha_n \mathbf{f}_n$$

Conducting Plate

Matrix equation:

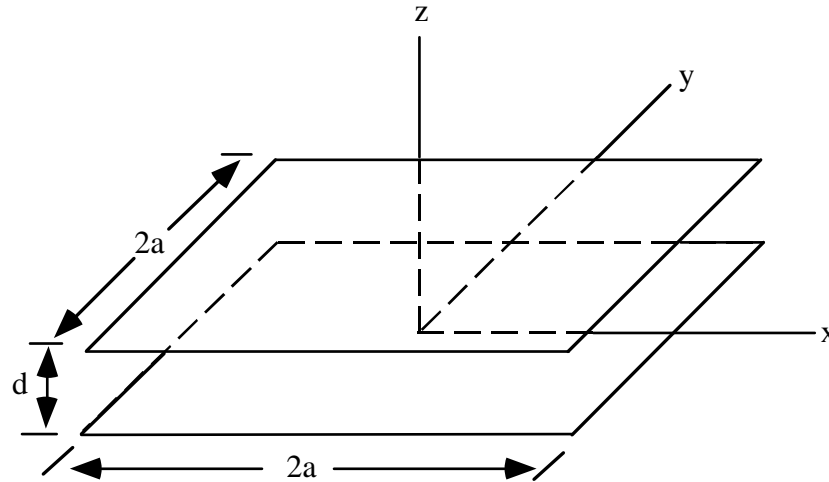
$$V = \sum_{n=1}^N l_{mn} f_n$$

Matrix element:

$$l_{mn} = \int_{\Delta x_m} dx' \int_{\Delta y_n} dy' \frac{1}{4\pi\epsilon \sqrt{(x_m - x')^2 + (y_n - x')^2}}$$

$$C = \frac{1}{V} \sum_{n=1}^N \alpha_n \Delta s_n = \sum_{mn} l_{mn}^{-1} \Delta s_n$$

Parallel Plates



Using N unknowns per plate, we get $2N \times 2N$ matrix equation:

$$[l] = \begin{bmatrix} [l^{tt}] & [l^{tb}] \\ [l^{bt}] & [l^{bb}] \end{bmatrix}$$

Subscript 't' for top and 'b' for bottom plate, respectively.

Parallel Plates

Matrix equation becomes $\begin{bmatrix} l_{mn}^{tt} & -l_{mn}^{tb} \end{bmatrix} \begin{bmatrix} \alpha_n^t \end{bmatrix} = \begin{bmatrix} g_m^t \end{bmatrix}$

Solution: $\begin{bmatrix} \alpha_m^t \end{bmatrix} = \begin{bmatrix} \left(l^{tt} - l^{tb} \right)_{mn}^{-1} \end{bmatrix} \begin{bmatrix} g_n^t \end{bmatrix}$

Capacitance $C = \frac{\text{charge on top plate}}{2V}$

$$= \frac{1}{2V} \sum_{top} \alpha_n^t \Delta s_n$$

Using $\Delta s = 4b^2$ and all elements of $\begin{bmatrix} g^t \end{bmatrix} = V$

$$C = 2b^2 \sum_{mn} \left(l^{tt} - l^{tb} \right)_{mn}^{-1}$$

Inductance

Relation: $\Psi = Li$

Ψ : flux stored by inductor

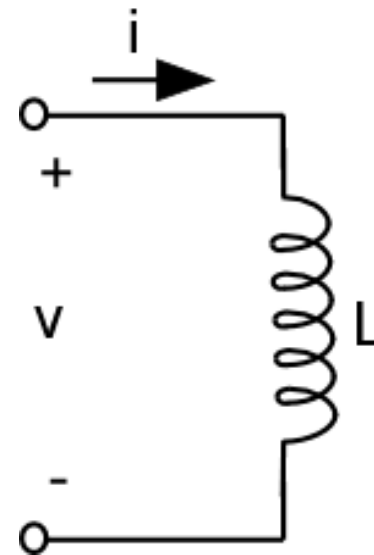
L : inductance

i : current through inductor

v : voltage across inductor

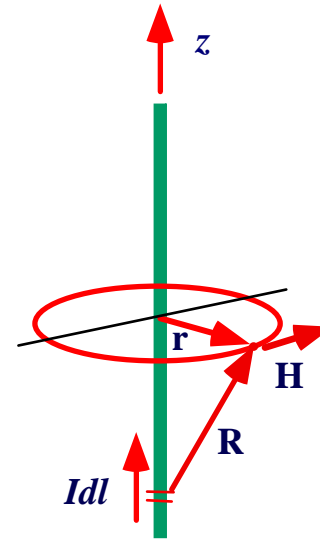
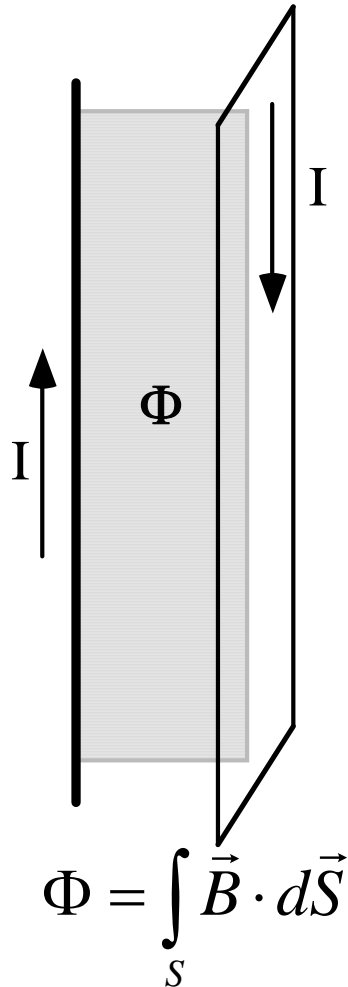
$$v(t) = L \frac{di}{dt} = \frac{d\Psi}{dt}$$

$$i(t) = \frac{1}{L} \int_0^t v(\tau) d\tau$$



Inductance

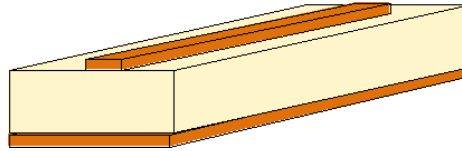
Magnetic Flux



$$\text{Inductance} = \frac{\text{Total flux linked}}{\text{Current}}$$

2-D Isomorphism

Electrostatics



Magnetostatics

$$(\hat{z} \times \hat{n}) \cdot \nabla_t V_i = 0$$

$$(\hat{z} \times \hat{n}) \cdot \nabla_t A_{zi} = 0$$

$$\hat{n} \cdot (\epsilon_{ri} \nabla_t V_i) = -\frac{q_s}{\epsilon_o}$$

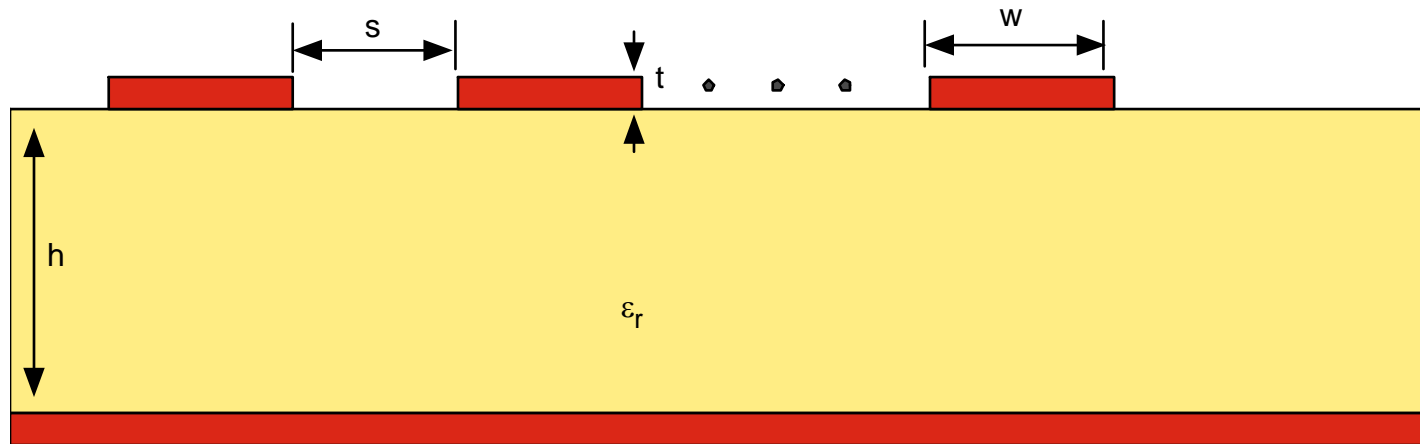
$$\hat{n} \cdot \left(\frac{1}{\mu_{ri}} \nabla_t A_{zi} \right) = -\mu_o J_z$$

$$CV = Q$$

$$LI = \psi$$

Consequence: 2D inductance can be calculated from 2D capacitance formulas

2-D N-line LC Extractor using MOM



- Symmetric signal traces
- Uniform spacing
- Lossless lines
- Uses MOM for solution

Output from MoM Extractor

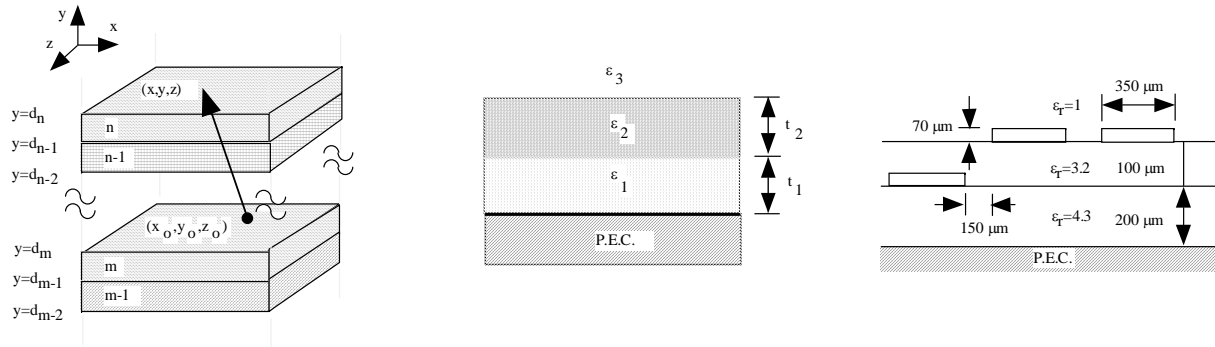
Capacitance (pF/m)

118.02299	-8.86533	-0.03030	-0.00011	-0.00000
-8.86533	119.04875	-8.86185	-0.03029	-0.00011
-0.03030	-8.86185	119.04876	-8.86185	-0.03030
-0.00011	-0.03029	-8.86185	119.04875	-8.86533
-0.00000	-0.00011	-0.03030	-8.86533	118.02299

Inductance (nH/m)

312.71680	23.42397	1.83394	0.14361	0.01128
23.42397	311.76042	23.34917	1.82812	0.14361
1.83394	23.34917	311.75461	23.34917	1.83394
0.14361	1.82812	23.34917	311.76042	23.42397
0.01128	0.14361	1.83394	23.42397	312.71680

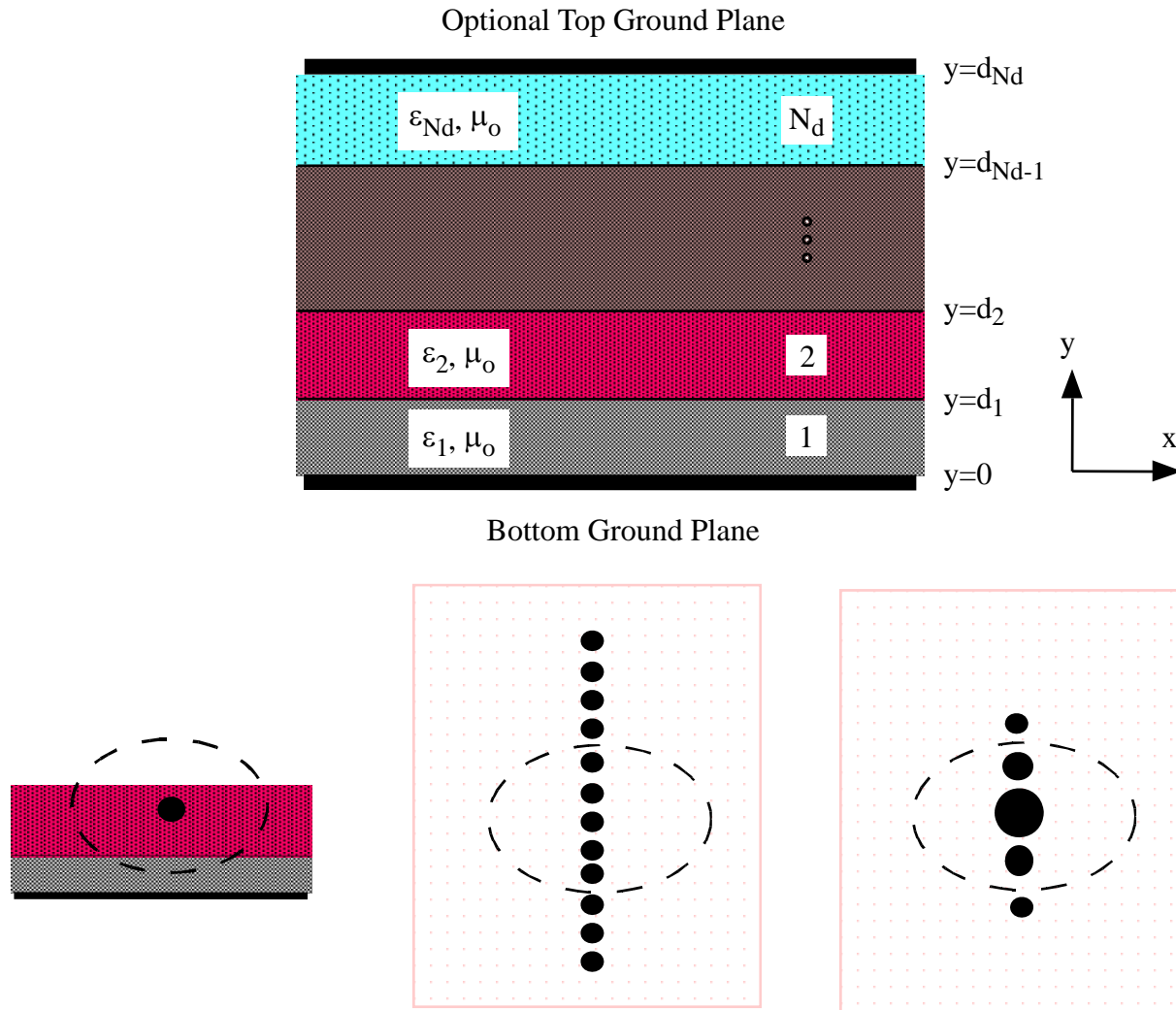
RLGC: Formulation Method



Closed-Form Spatial Green's Function

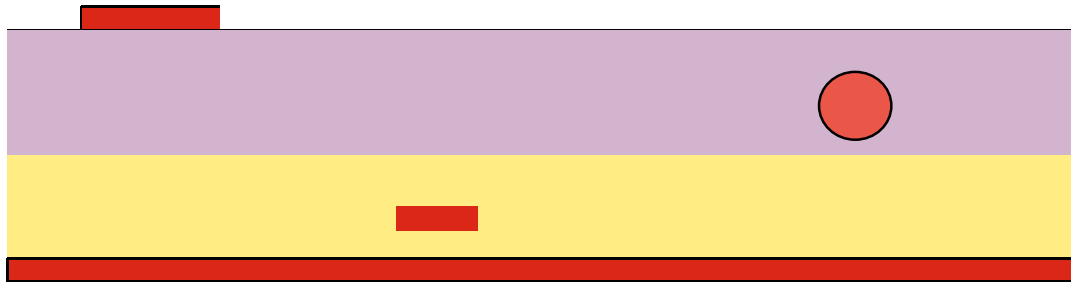
- * Computes 2-D and 3-D capacitance matrix in multilayered dielectric
- * Method is applicable to arbitrary polygon-shaped conductors
- * Computationally efficient
- **Reference**
 - K. S. Oh, D. B. Kuznetsov and J. E. Schutt-Aine, "Capacitance Computations in a Multilayered Dielectric Medium Using Closed-Form Spatial Green's Functions," IEEE Trans. Microwave Theory Tech., vol. MTT-42, pp. 1443-1453, August 1994.

Multilayer Green's Function



Extraction Program: RLGC

RLGC computes the four transmission line parameters, viz., the capacitance matrix **C**, the inductance matrix **L**, the conductance matrix **G**, and the resistance matrix **R**, of a multiconductor transmission line in a multilayered dielectric medium. **RLGC** features the following list of functions:



RLGC – Multilayer Extractor

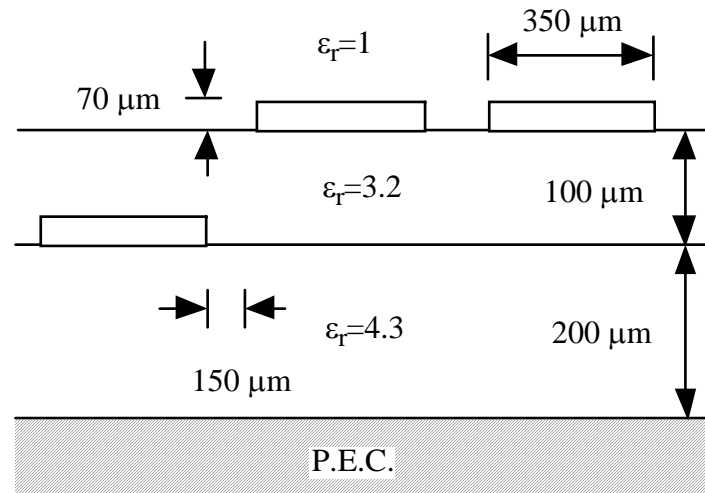
- **Features**

- Handling of dielectric layers with no ground plane, either top or bottom ground plane (microstrip cases), or both top and bottom ground planes (stripline cases)
- Static solutions are obtained using the Method of Moment (MoM) in conjunction with the recently-developed closed-form Green's functions: one of the most accurate and efficient methods for static analysis
- Modeling of dielectric losses as well as conductor losses (including ground plane losses)
- The resistance matrix R is computed based on the current distribution - more accurate than the use of any closed-form formulae
- Both the proximity effect and the skin effect are modeled in the resistance matrix R .
- Computes the potential distribution
- Handling of an arbitrary number of dielectric layers as well as an arbitrary number of conductors.
- The cross section of a conductor can be arbitrary or even be infinitely thin

- **Reference**

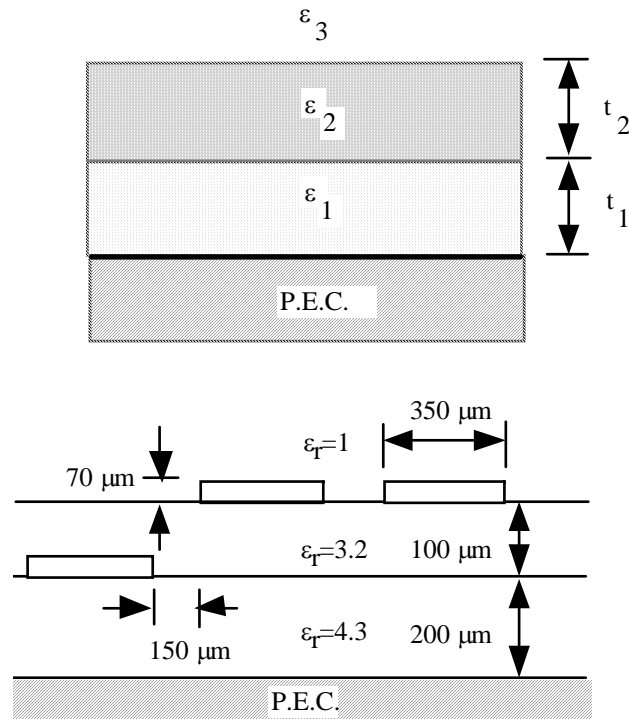
- K. S. Oh, D. B. Kuznetsov and J. E. Schutt-Aine, "Capacitance Computations in a Multilayered Dielectric Medium Using Closed-Form Spatial Green's Functions," IEEE Trans. Microwave Theory Tech., vol. MTT-42, pp. 1443-1453, August 1994.

RLGC – General Topology



Three conductors in a layered medium. All conductor dimensions and spacing are identical. The loss tangents of the lower and upper dielectric layers are 0.004 and 0.001 respectively, the conductivity of each line is $5.8e7$ S/m, and the operating frequency is 1 GHz

3-Line Capacitance Results



Capacitance Matrix (pF/m)

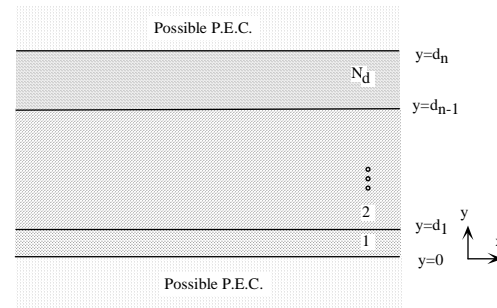
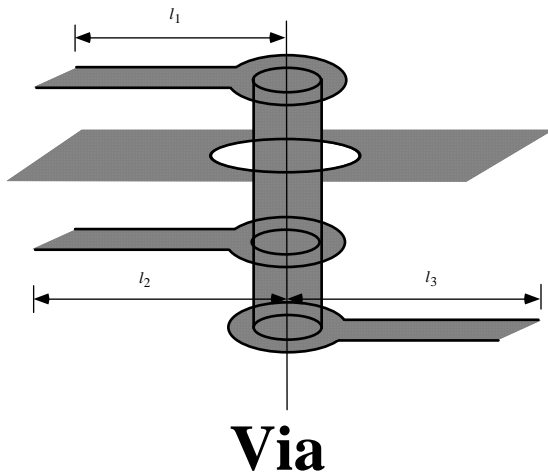
$$\begin{bmatrix} 142.09 & -21.765 & -0.8920 \\ -21.733 & 93.529 & -18.098 \\ -0.8900 & -18.097 & 87.962 \end{bmatrix}$$

Delabare et al.

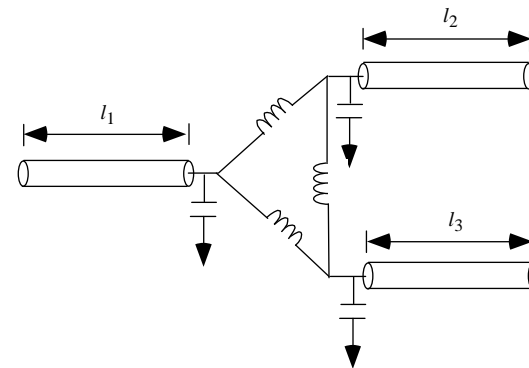
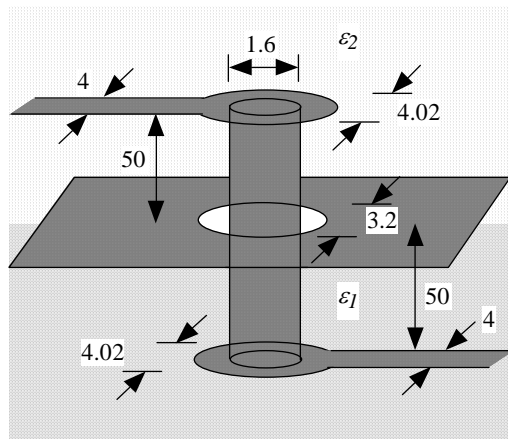
$$\begin{bmatrix} 145.33 & -23.630 & -1.4124 \\ -22.512 & 93.774 & -17.870 \\ -1.3244 & -17.876 & 87.876 \end{bmatrix}$$

RLGC Method

Modeling Vias

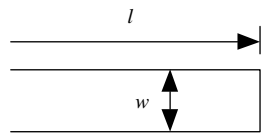


Medium

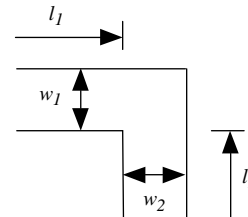
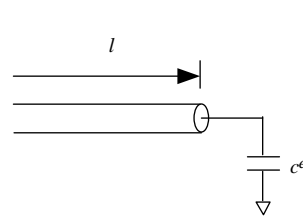


Equivalent circuit

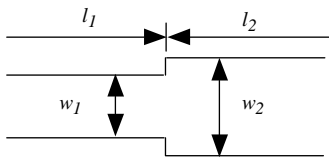
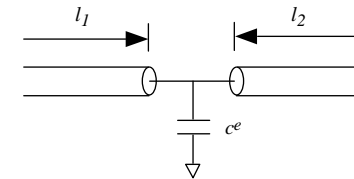
Modeling Discontinuities



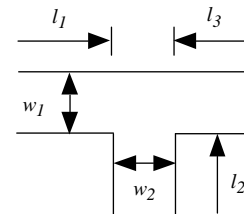
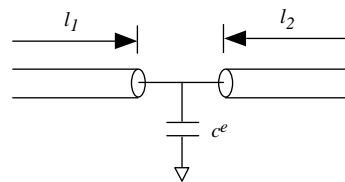
Open



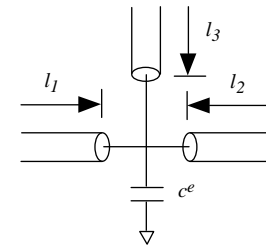
Bend



Step

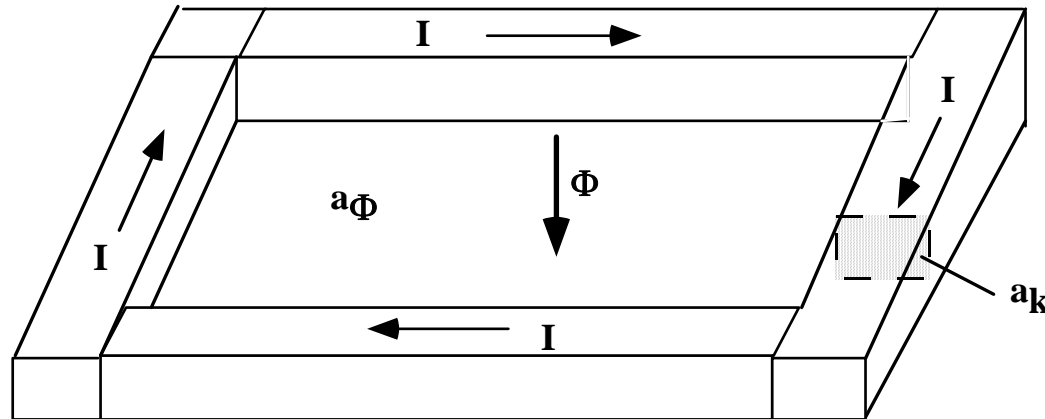


T-Junction



3D Inductance Calculation

Loop Inductance

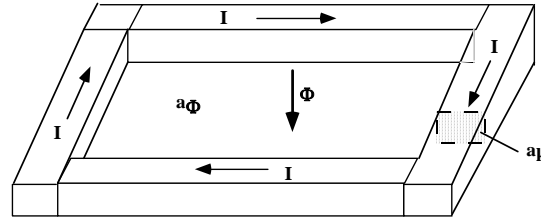


$$L_{loop} = \frac{\Phi}{I} = \frac{1}{I} \int_{a_\Phi} \vec{B} \cdot d\vec{a} = \frac{1}{I} \int_{a_\Phi} (\nabla \times \vec{A}) \cdot d\vec{a}$$

QUESTION: Can we associate inductance with piece of conductor rather than a loop? \rightarrow PEEC Method

Partial Inductance (PEEC) Approach

QUESTION: Can we associate inductance with piece of conductor rather than a loop?



$$L_{loop} = \sum_{i=1}^4 \sum_{j=1}^4 \frac{1}{a_i a_j} \int_{a_i} \int_{a_j} \int_{l_i} \int_{l_j} \frac{\mu}{4\pi} \frac{d\vec{l}_i \cdot d\vec{l}_j}{|\vec{r}_i - \vec{r}_j|} da_i da_j$$

DEFINITION OF PARTIAL INDUCTANCE

$$L_{pij} = \frac{1}{a_i a_j} \frac{\mu}{4\pi} \int_{a_i} \int_{a_j} \int_{l_i} \int_{l_j} \frac{d\vec{l}_i \cdot d\vec{l}_j}{|\vec{r}_i - \vec{r}_j|} da_i da_j$$

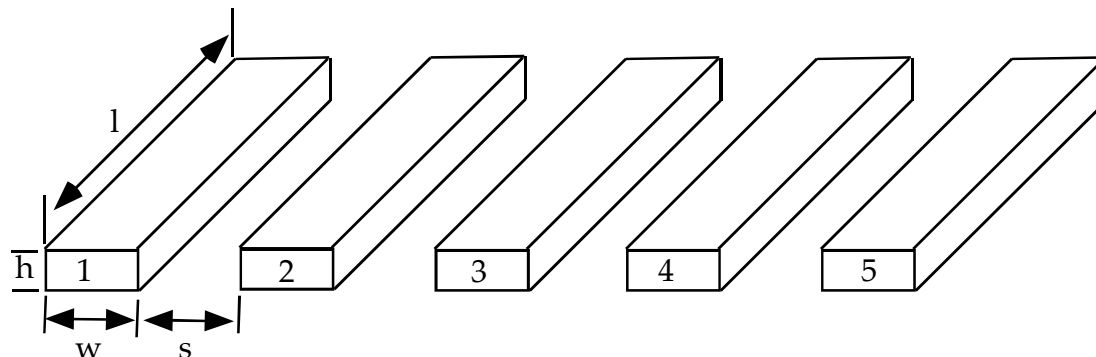
$$L_{loop} = \sum_{i=1}^4 \sum_{j=1}^4 s_{ij} L_{pij}$$

Circuit Element K

$$[K]=[L]^{-1}$$

- Better locality property
- Leads to sparser matrix
- Diagonally dominant
- Allows truncation of far off-diagonal elements
- Better suited for on-chip inductance analysis

Locality of K Matrix



$$[L] = \begin{bmatrix} 11.4 & 4.26 & 2.54 & 1.79 & 1.38 \\ 4.26 & 11.4 & 4.26 & 2.54 & 1.79 \\ 2.54 & 4.26 & 11.4 & 4.26 & 2.54 \\ 1.79 & 2.54 & 4.26 & 11.4 & 4.26 \\ 1.38 & 1.79 & 2.54 & 4.26 & 11.4 \end{bmatrix}$$

$$[K] = \begin{bmatrix} 103 & -34.1 & -7.80 & -4.31 & -3.76 \\ -34.1 & 114 & -31.6 & -6.67 & -4.31 \\ -7.80 & -31.6 & 115 & -31.6 & -7.80 \\ -4.31 & -6.67 & -31.6 & 114 & -34.1 \\ -3.76 & -4.31 & -7.80 & -34.1 & 103 \end{bmatrix}$$

Package Inductance & Capacitance

<u>Component</u>	<u>Capacitance</u> (pF)	<u>Inductance</u> (nH)
68 pin plastic DIP pin [†]	4	35
68 pin ceramic DIP pin ^{††}	7	20
68 pin SMT chip carrier [†]	2	7

† No ground plane; capacitance is dominated by wire-to-wire component.

†† With ground plane; capacitance and inductance are determined by the distance between the lead frame and the ground plane, and the lead length.

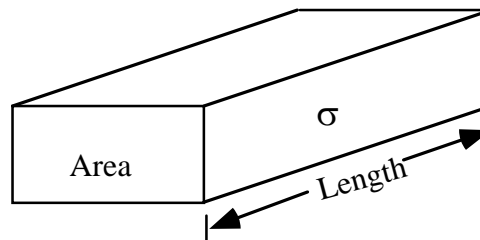
Package Inductance & Capacitance

<u>Component</u>	<u>Capacitance</u> (pF)	<u>Inductance</u> (nH)
68 pin PGA pin††	2	7
256 pin PGA pin††	5	15
Wire bond	1	1
Solder bump	0.5	0.1

† No ground plane; capacitance is dominated by wire-to-wire component.

†† With ground plane; capacitance and inductance are determined by the distance between the lead frame and the ground plane, and the lead length.

Metallic Conductors



Resistance : R

$$R = \frac{\text{Length}}{\sigma \cdot \text{Area}}$$

Package level:

W=3 mils

R=0.0045 Ω /mm

Submicron level:

W=0.25 microns

R=422 Ω /mm

Metallic Conductors

Metal

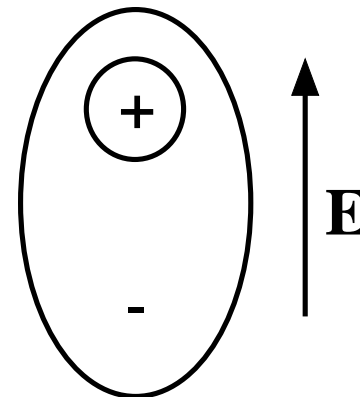
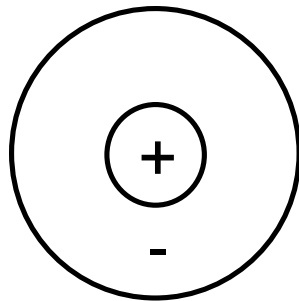
Conductivity

σ ($\Omega^{-1} \text{ m}^{-1} \times 10^{-7}$)

Silver	6.1
Copper	5.8
Gold	3.5
Aluminum	1.8
Tungsten	1.8
Brass	1.5
Solder	0.7
Lead	0.5
Mercury	0.1

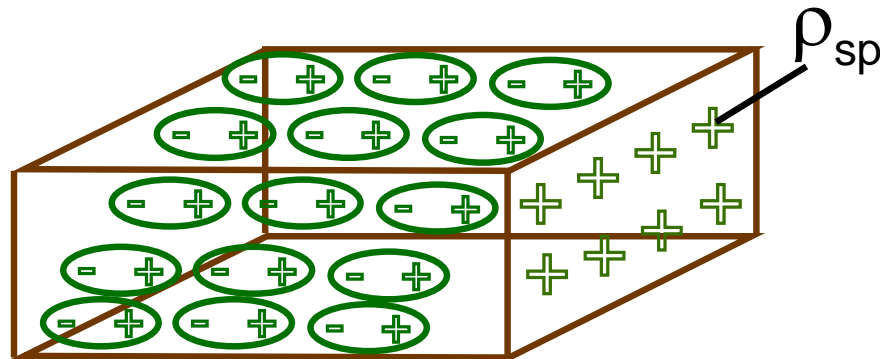
Dielectrics

- Dielectrics contain charges that are tightly bound to the nuclei
- Charges can move a fraction of an atomic distance away from equilibrium position
- Electron orbits can be distorted when an electric field is applied



Dielectrics

- Charge density within volume is zero
- Surface charge density is nonzero



$$\mathbf{D} = \epsilon_0(1 + \chi_e)\mathbf{E} = \epsilon\mathbf{E}$$

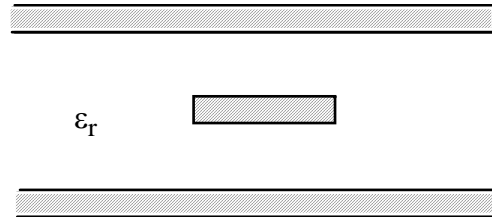
Dielectric Materials

$$v = \sqrt{\frac{1}{LC}}$$

<u>Material</u>	<u>Conductivity</u> ($\Omega^{-1}\text{-m}^{-1}$)
Germanium	2.2
Silicon	0.0016
Glass	10^{-10} - 10^{-14}
Quartz	0.5×10^{-17}

$$\tan \delta = \frac{\sigma}{\omega \epsilon}$$

Dielectric Materials



$$v = \sqrt{\frac{1}{LC}}$$

<u>Material</u>	<u>ϵ_r</u>	<u>v(cm/s)</u>
Polyimide	2.5-3.5	16-19
Silicon dioxide	3.9	15
Epoxy glass (FR4)	5.0	13
Alumina (ceramic)	9.5	10

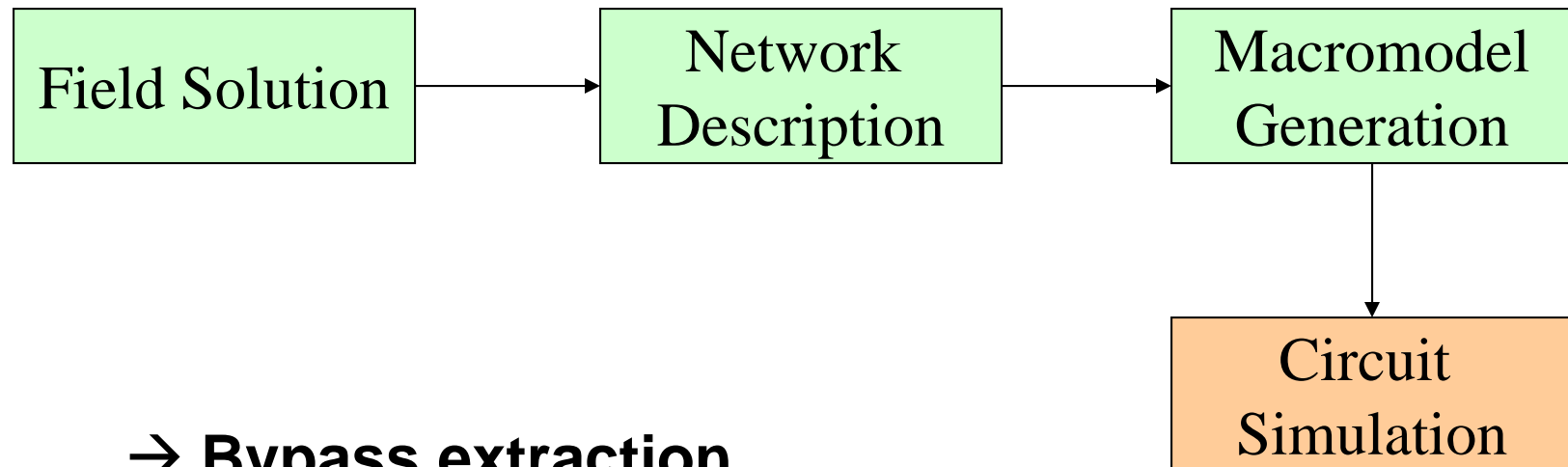
Conductivity of Dielectric Materials

$$\epsilon = \epsilon_r + j \epsilon_i$$

<u>Material</u>	<u>Conductivity ($\Omega^{-1} \text{ m}^{-1}$)</u>
Germanium	2.2
Silicon	0.0016
Glass	$10^{-10} - 10^{-14}$
Quartz	0.5×10^{-17}

$$\text{Loss TANGENT : } \tan\delta = \frac{\sigma}{\omega \epsilon}$$

Combining Field and Circuit Solutions



→ Bypass extraction procedure through the use of Y, Z, or S parameters (frequency domain)

Full-Wave Methods

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \text{Faraday's Law of Induction}$$

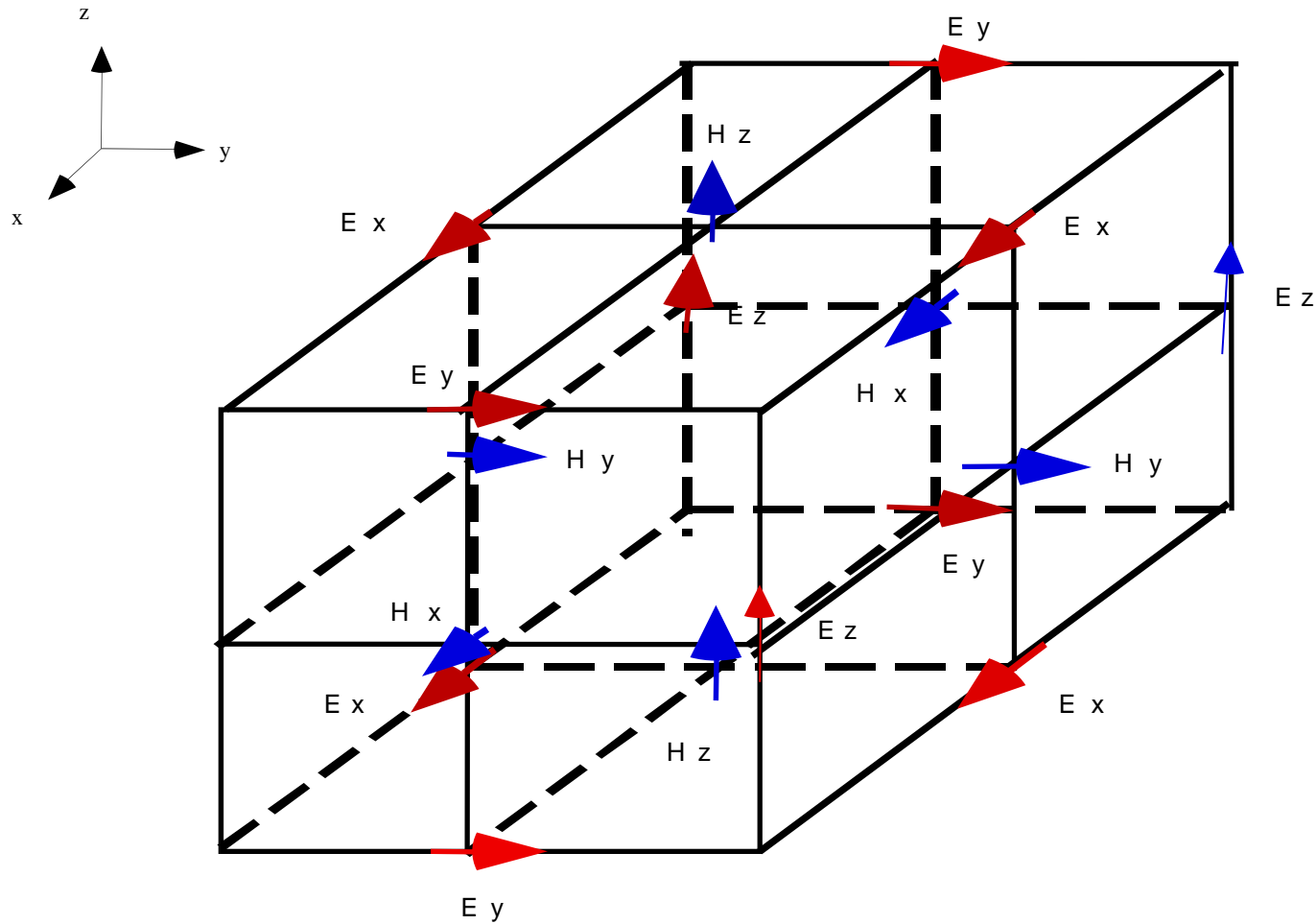
$$\nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t} \quad \text{Ampère's Law}$$

$$\nabla \cdot \vec{D} = \rho \quad \text{Gauss' Law for electric field}$$

$$\nabla \cdot \vec{B} = 0 \quad \text{Gauss' Law for magnetic field}$$

FDTD: Discretize equations and solve with appropriate boundary conditions

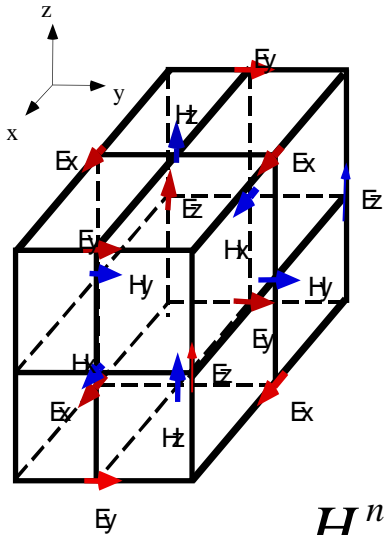
Finite Difference Time Domain (FDTD)



Space Discretization

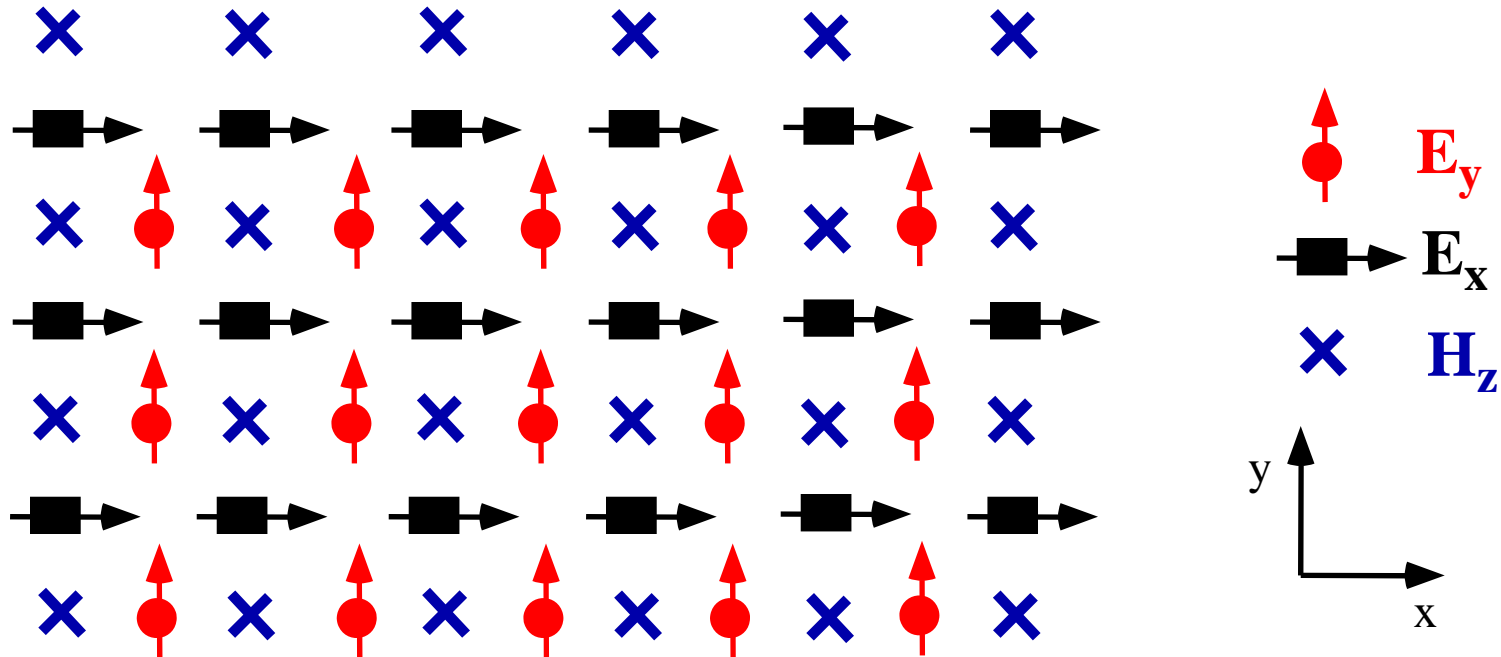
FDTD – Yee Algorithm

$$E_x^n(i, j, k) = E_x^{n-1} + \frac{c}{\epsilon} \frac{\Delta t}{\Delta y} \left(H_z^{n-1/2}(i, j, k) - H_z^{n-1/2}(i, j-1, k) \right) - \frac{c}{\epsilon} \frac{\Delta t}{\Delta z} \left(H_y^{n-1/2}(i, j, k) - H_y^{n-1/2}(i, j, k-1) \right)$$



$$H_x^{n+1/2}(i, j, k) = H_x^{n-1/2} + \frac{c}{\mu} \frac{\Delta t}{\Delta y} \left(E_z^n(i, j+1, k) - E_z^n(i, j, k) \right) + \frac{c}{\mu} \frac{\Delta t}{\Delta z} \left(E_y^{n-1/2}(i, j, k+1) - E_y^n(i, j, k) \right)$$

2D-FDTD



$$E_x^n\left(i + \frac{1}{2}, j\right) = E_x^{n-1}\left(i + \frac{1}{2}, j\right) + \frac{\Delta t}{\epsilon_0 \Delta y} \left[H_z^{n-1/2}\left(i + \frac{1}{2}, j + \frac{1}{2}\right) - H_z^{n-1/2}\left(i + \frac{1}{2}, j - \frac{1}{2}\right) \right]$$

$$E_y^n\left(i, j + \frac{1}{2}\right) = E_y^{n-1}\left(i, j + \frac{1}{2}\right) - \frac{\Delta t}{\epsilon_0 \Delta x} \left[H_z^{n-1/2}\left(i + \frac{1}{2}, j + \frac{1}{2}\right) - H_z^{n-1/2}\left(i - \frac{1}{2}, j + \frac{1}{2}\right) \right]$$

$$H_z^{n+1/2}\left(i + \frac{1}{2}, j + \frac{1}{2}\right) = H_z^{n-1/2}\left(i + \frac{1}{2}, j + \frac{1}{2}\right) + \frac{\Delta t}{\mu_0 \Delta y} \left[E_x^n\left(i + \frac{1}{2}, j + 1\right) - E_x^n\left(i + \frac{1}{2}, j\right) \right] \\ - \frac{\Delta t}{\mu_0 \Delta x} \left[E_y^n\left(i + 1, j + \frac{1}{2}\right) - E_y^n\left(i, j + \frac{1}{2}\right) \right]$$

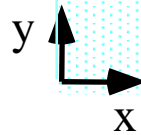
Absorbing Boundary Condition: 2D-PML Formulation

Simulation Medium

$$\begin{aligned}\epsilon_0 \frac{\partial E_x}{\partial t} &= \frac{\partial H_z}{\partial y} \\ \epsilon_0 \frac{\partial E_y}{\partial t} &= -\frac{\partial H_z}{\partial x} \\ \mu_0 \frac{\partial H_z}{\partial t} &= \frac{\partial E_x}{\partial y} - \frac{\partial E_y}{\partial x}\end{aligned}$$

PML Medium

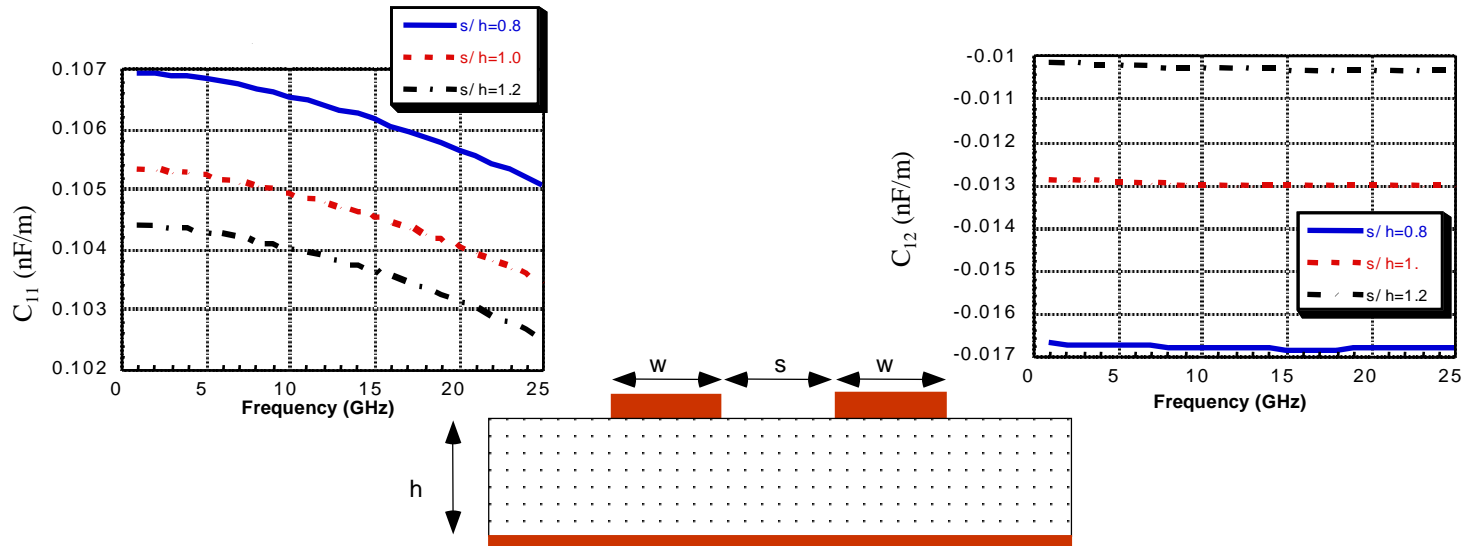
$$\begin{aligned}\epsilon_0 \frac{\partial E_x}{\partial t} + \sigma E_x &= \frac{\partial H_z}{\partial y} \\ \epsilon_0 \frac{\partial E_y}{\partial t} + \sigma E_y &= -\frac{\partial H_z}{\partial x} \\ \mu_0 \frac{\partial H_z}{\partial t} + \sigma^* H_z &= \frac{\partial E_x}{\partial y} - \frac{\partial E_y}{\partial x}\end{aligned}$$



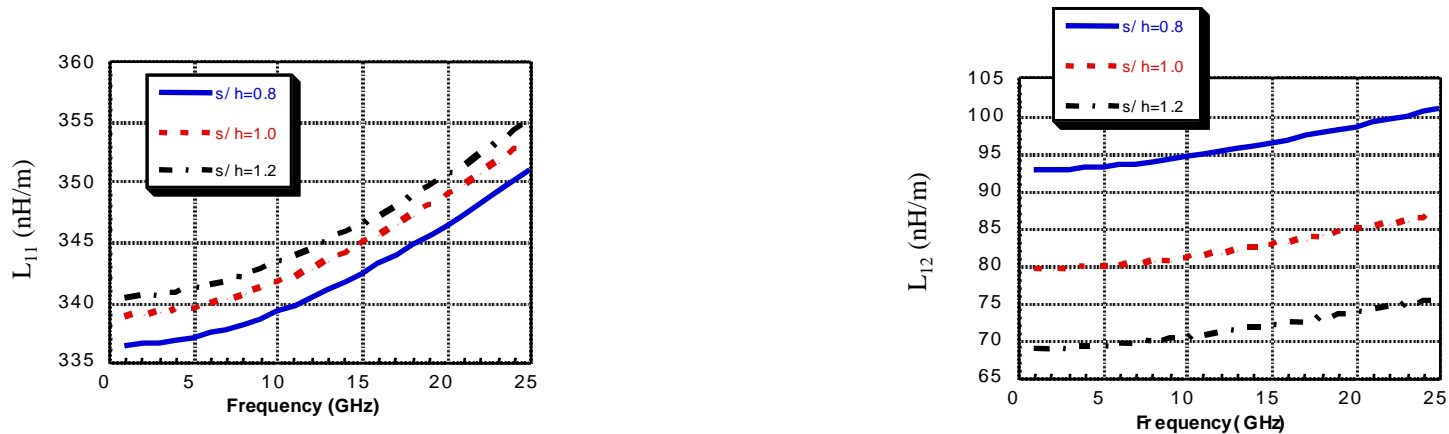
$$\frac{\sigma}{\epsilon_0} = \frac{\sigma^*}{\mu_0}$$

No reflection from PML interface

FDTD – Coupled Microstrips



Parameters: $w=0.3\text{mm}$, $h=0.25\text{mm}$. Dielectric constant is 4.5.



Parallel-FDTD - Motivations

PCB interconnect EM problems

1. Interconnect coupling / crosstalk
2. Signal distortion
3. Radiation (EMI)
4. etc

PCB modeling challenges:

1. Size of the problem - determined by the ratio of board size and finest features
2. Simulations time

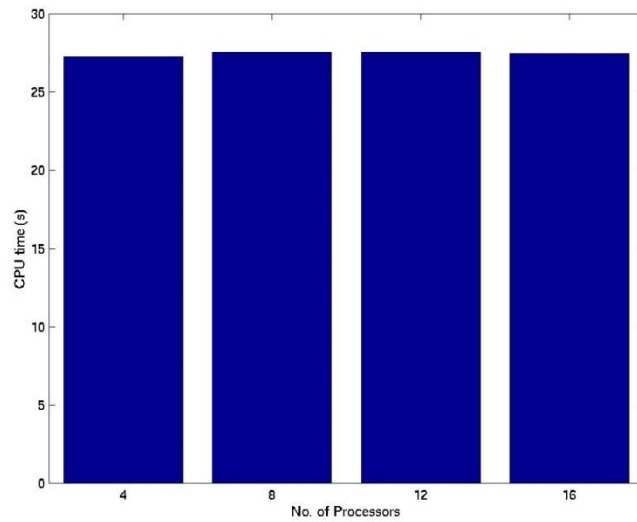
Solutions: ~~Supercomputer~~

Mini-computer cluster

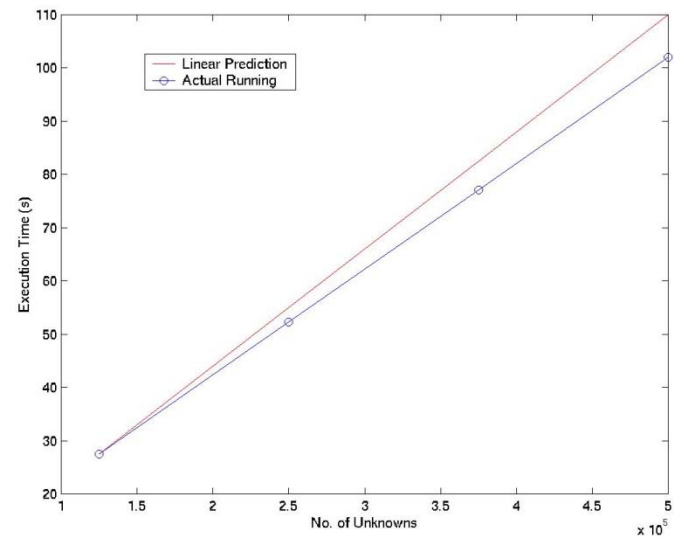
Gflops/dollar

Parallel-FDTD

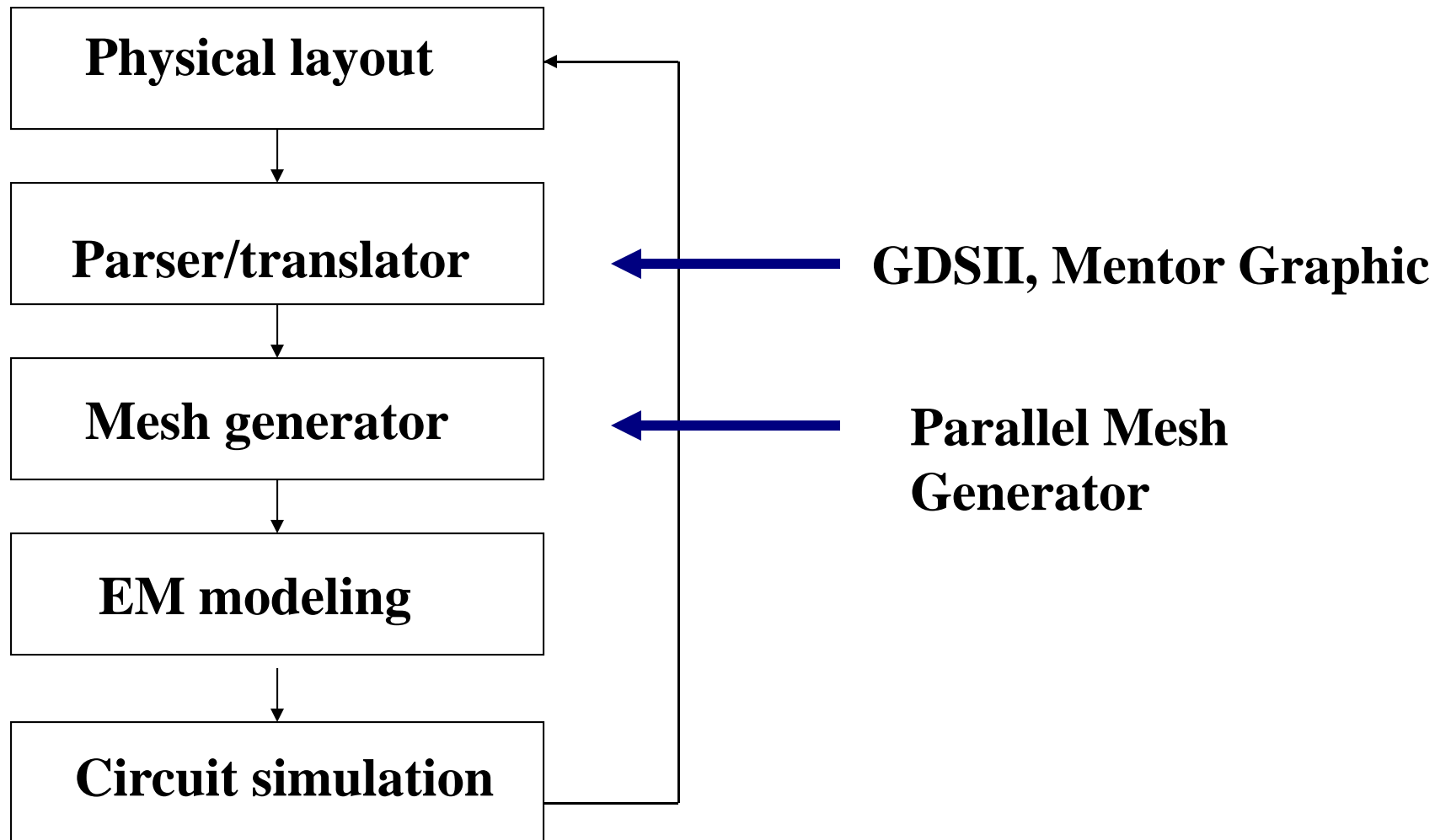
Scalability



Speed vs. problem size



Parallel-FDTD



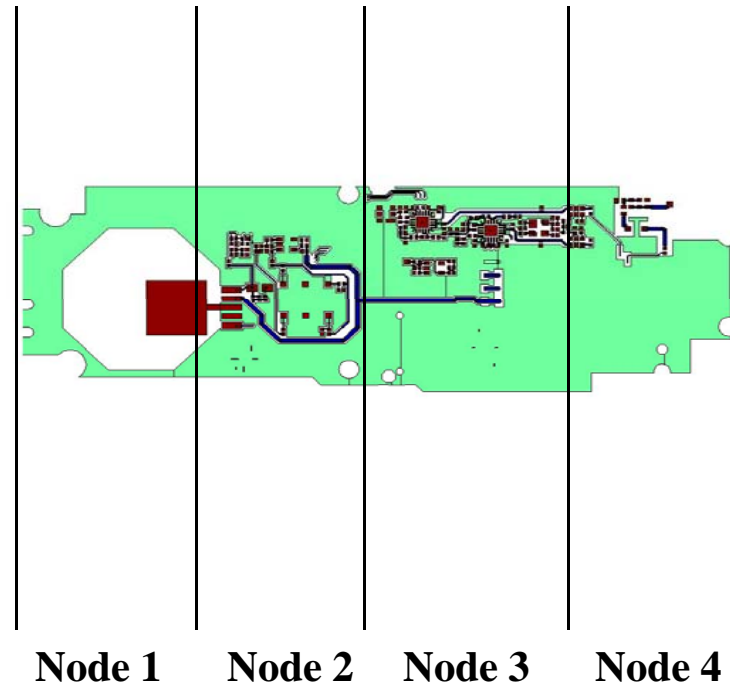
Parallel-FDTD

Max(computation/communication)

Maximize the ratio of volume to surface



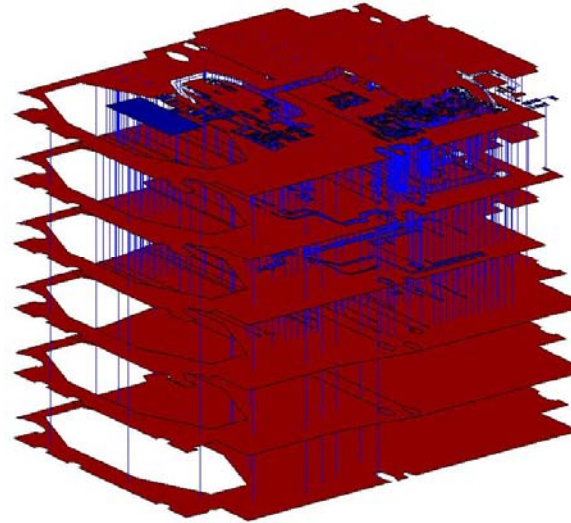
Partition along the largest dimension



Parallel-FDTD

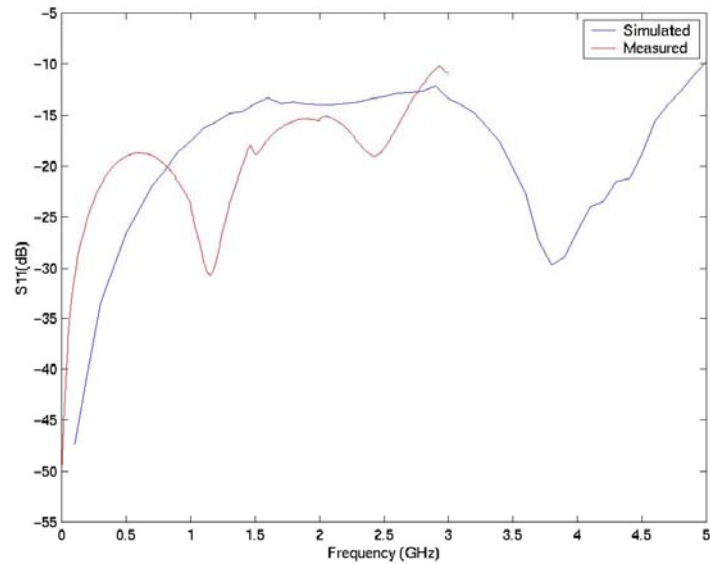
A cellular phone PCB test board

- **six layered test board**
- **over 4000 vias**
- **dielectric materials FR 4**
- **total thickness < 20 mils**
- **over 60 traces**
- **around 30 PEC fills**
- **around 300 component pads**



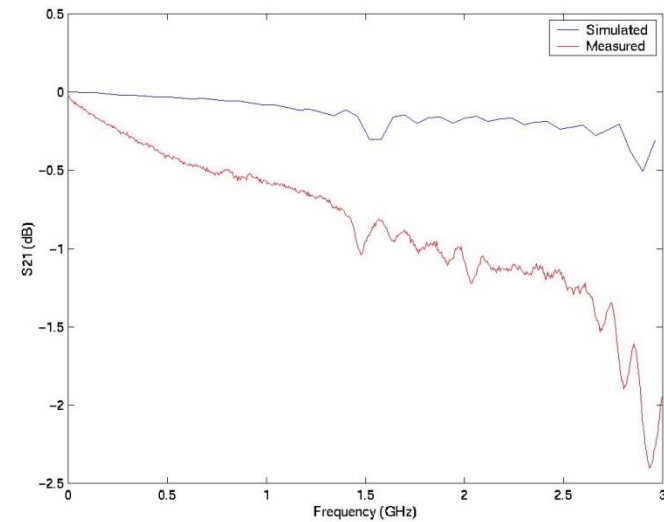
Parallel-FDTD

Comparison between measurement and simulation



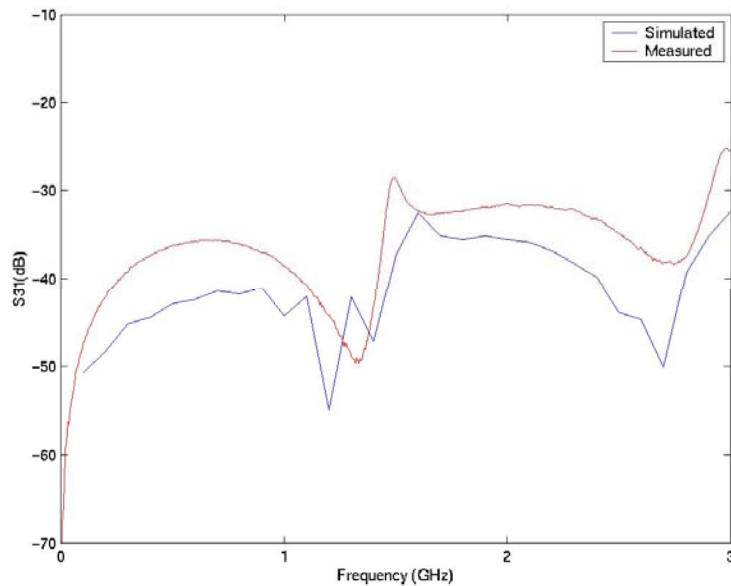
S11

S21

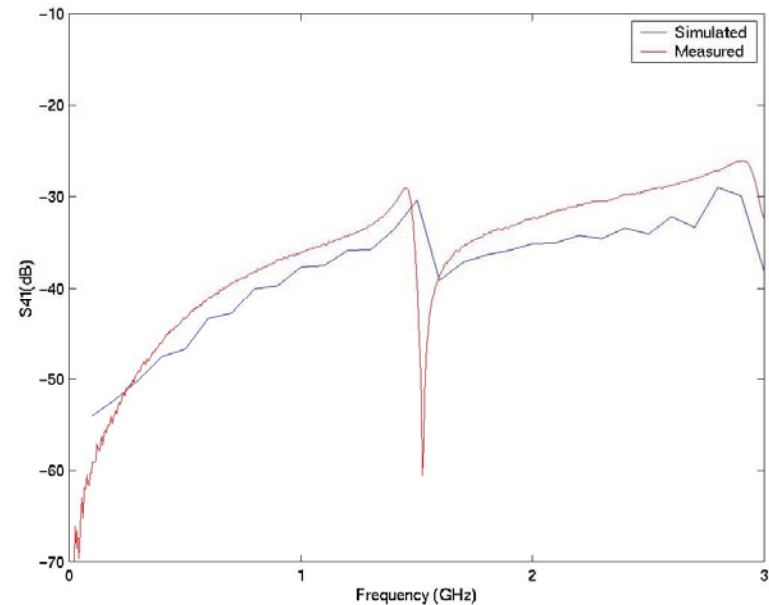


Parallel-FDTD

S31



S41



- **Reference**

- F. Liu, J. Schutt-Ainé and J. Chen, "Full Wave Analysis and Modeling of Multiconductor Transmission Lines via 2-D-FDTD and Signal-Processing Techniques," IEEE Trans. Microwave Theory Tech. vol. 50, pp. 570-577, February 2002.